

## Problem set 1

1. Consider a 1D lattice with lattice constant  $a$ , and a particle whose dynamics is described by the hopping Hamiltonian  $\hat{H}_0 = -t \sum_n (|n\rangle\langle n+1| + |n+1\rangle\langle n|)$ .

(a) Find  $G_0(n_2, n_1, E)$ . Suggested steps: (i) assume a finite number of lattice sites,  $N$ , and periodic boundary conditions (i.e, site " $N+1$ " is in fact site 1). Find the eigenenergies and eigenfunctions. (ii) write  $G_0(n_2, n_1, E)$  in terms of eigenstates and eigenenergies. Take the limit  $N \rightarrow \infty$ , and evaluate the resulting integral. It's fine with me if you pick the result from a table of integrals, though you should be able to do it by hand.

(b) If  $E$  is close to the bottom of the band, the dispersion energy can be approximated (up to a constant) with a quadratic form similar to the energy of a free particle. Find the corresponding effective mass. In this limit ( $a \rightarrow 0$ ), is  $G_0(n_2, n_1, E)$  in agreement with the Green's function for a free 1D particle discussed in class?

2. Consider a spin- $\frac{1}{2}$  whose dynamics is described by the Hamiltonian  $\hat{H} = -b\hat{S}_x$ . Find (a)  $G(\uparrow, \downarrow; E)$ ; (b) Fourier transform this to obtain  $G(\uparrow t; \downarrow 0) = -i/\hbar \mathcal{G}(\uparrow t; \downarrow 0)$ . Now check your result by solving directly the Schrödinger equation for a particle which has spin  $\downarrow$  at  $t = 0$ , and finding the amplitude of probability to have spin  $\uparrow$  at a later time  $t$ .

3. Consider an electron in a sample of volume  $V$ . The electron interacts with a set of  $N$  impurities located at random positions  $\vec{R}_1, \dots, \vec{R}_N$ , through the potential  $V_T(\vec{r}) = \sum_{n=1}^N V(\vec{r} - \vec{R}_n)$ . We assume that there are very many impurities, but each one is very weak. Also, we assume that  $\int d\vec{r} V(\vec{r}) = 0$ .

Consider now disorder averages. By definition

$$\langle f(\vec{R}_1, \dots, \vec{R}_N) \rangle_{dis} = \int \frac{d\vec{R}_1}{V} \dots \int \frac{d\vec{R}_N}{V} f(\vec{R}_1, \dots, \vec{R}_N)$$

(a) calculate  $\langle V_T(\vec{r}_1) V_T(\vec{r}_2) \rangle = W(\vec{r}_1 - \vec{r}_2)$ . Express this in terms of the Fourier components of the impurity potential  $V_{\vec{q}} = V_{-\vec{q}}^* = \int d\vec{r} e^{-i\vec{q}\vec{r}} V(\vec{r})$ .

(b) calculate  $\langle V_T(\vec{r}_1) V_T(\vec{r}_2) V_T(\vec{r}_3) V_T(\vec{r}_4) \rangle$ . Does it satisfy the factorization rule of a gaussian disorder? If not, argue in what case could we still assume that factorization to be a good approximation? Is this approximation likely to become better or worse, as you go to higher order correlations?

4. Consider an electron in an inhomogeneous external magnetic field,  $\vec{B}(\vec{r})$ , such that its Hamiltonian is  $\mathcal{H} = \mathcal{H}_0 - \vec{b}(\vec{r}) \cdot \hat{\vec{\sigma}}$ . Here,  $\mathcal{H}_0 = \hat{p}^2/2m$  is just the kinetic energy of a free electron, we expressed the spin electron in terms of Pauli matrices,  $\hat{\vec{s}} = \hbar/2\hat{\vec{\sigma}}$ , and I included all other constants (Bohr magneton, factor of  $\hbar$  etc), in rescaling  $\vec{b}(\vec{r}) \sim \vec{B}(\vec{r})$ . Use appropriate diagrams and rules to express Dyson's equation for this problem, in real space. Write explicitly the expression of the second-order contribution,  $G^{(2)}(\vec{r}, \sigma; \vec{r}', \sigma'; E)$ . Now, assume that  $\vec{h}(\vec{r}) = h\vec{e}_x$  (i.e., uniform field in the x-direction). Switch to  $\vec{k}$ -space, derive the corresponding diagrams and needed rules, and use Dyson's equation to find  $G_{\sigma\sigma'}(\vec{k}, E)$  exactly. Is the answer sensible?