

The Henon-Heiles Hamiltonian

Motion of Stars about a Galactic
Center

The Henon-Heiles Hamiltonian describes the motion of stars around a galactic center, assuming the motion is restricted to the xy plane.



$$H = \frac{P_x^2}{2m} + \frac{P_y^2}{2m} + \frac{1}{2}k(x^2 + y^2) + \lambda(x^2y - \frac{1}{3}y^3)$$

The Equations of Motion

In order to simplify the Hamiltonian assume:

$$P_x = m\dot{x} \quad \text{and} \quad P_y = m\dot{y}$$

And rewrite the Hamiltonian in a dimensionless, normalized form with lambda equal to one:

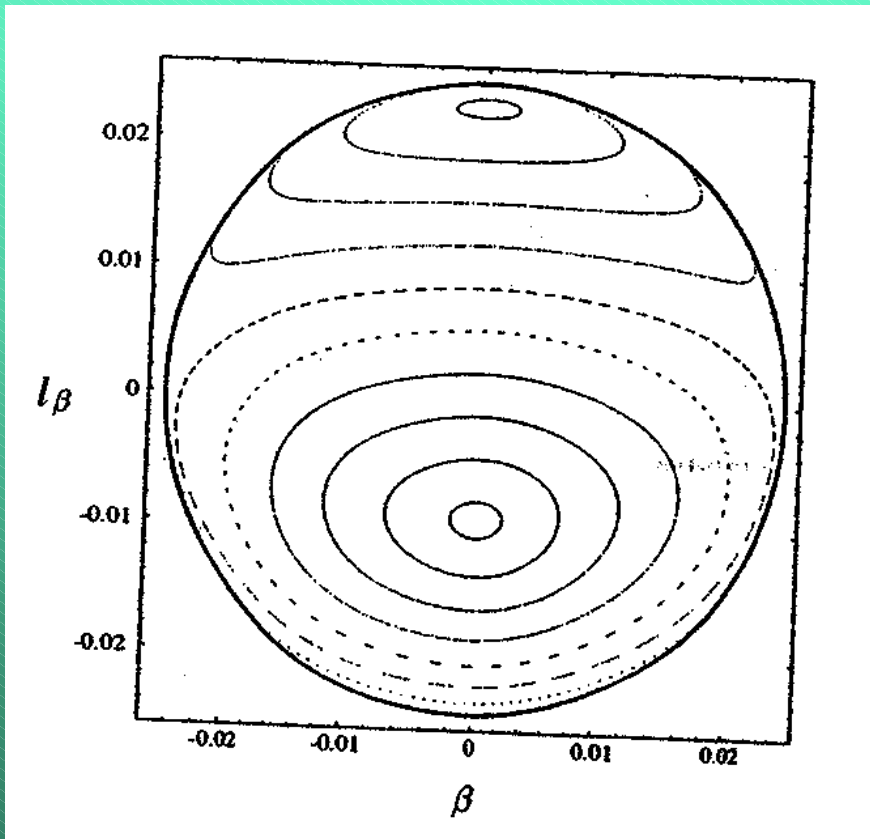
$$E = \frac{1}{2}(\dot{x}^2 + \dot{y}^2) + \frac{1}{2}(x^2 + y^2) + x^2y - \frac{1}{3}y^3$$

This yields the following equations of motion:

$$\ddot{x} = -x - 2xy \quad \text{and} \quad \ddot{y} = -y - x^2 + y^2$$

Poincare Maps

A Poincare map is a two dimensional slice of a systems four dimensional phase space



Poincare map for the double pendulum

Beta is the angle between the first and second pendulums

Step 1:

Determine Appropriate Initial Conditions

Appropriate initial conditions are determined by the total energy of the system

$$E = \frac{1}{2}(\dot{x}^2 + \dot{y}^2) + \frac{1}{2}(x^2 + y^2) + x^2y - \frac{1}{3}y^3$$

We will be examining Poincare sections in the $y\dot{y}$ plane, so initially we may set $x=0$

$$E = \frac{1}{2}(\dot{x}^2 + \dot{y}^2) + \frac{1}{2}y^2 - \frac{1}{3}y^3$$

The bounding curve for the available phase space can be determined by setting $x=0$

$$E = \frac{1}{2} \dot{y}^2 + \frac{1}{2} y^2 - \frac{1}{3} y^3$$

If we choose values of y and \dot{y} that lie within this bounding Curve, we may use the equation

$$\dot{x} = \left(2E - \dot{y}^2 - y^2 + \frac{2}{3} y^3 \right)^{\frac{1}{2}}$$

This gives us our appropriate set of initial conditions

Step 2:

Integrate the Equations of Motion

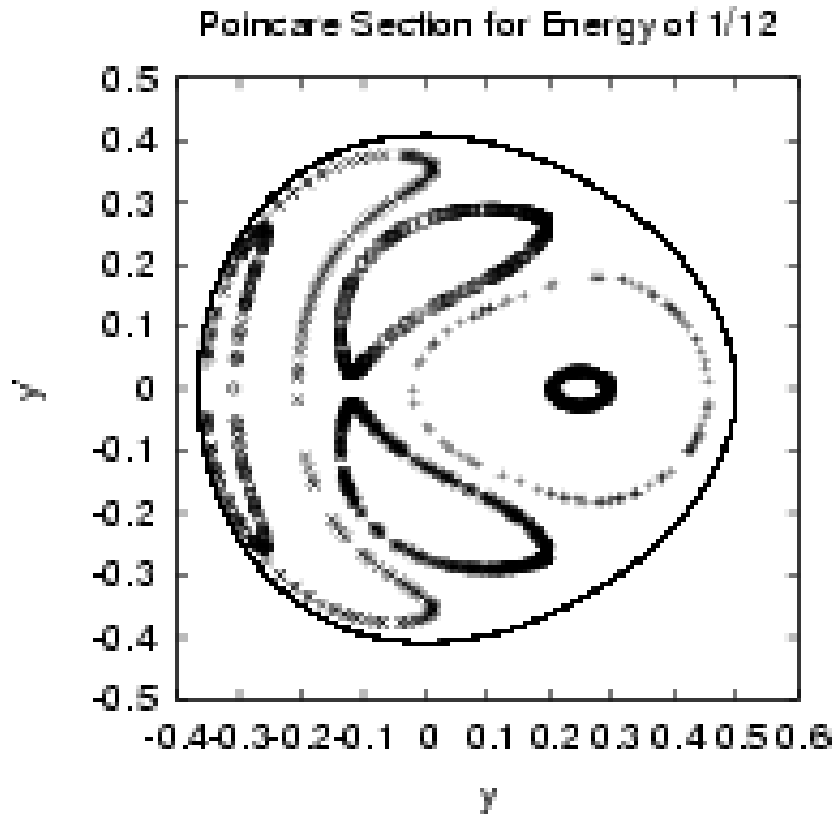
$$\ddot{x} = -x - 2xy \qquad \ddot{y} = -y - x^2 + y^2$$

In order to solve these equations we will first cast them into canonical, first order form

$$\begin{aligned} y_1 &= x & y_1' &= y_3 \\ y_2 &= y & y_2' &= y_4 \\ y_3 &= \dot{x} & y_3' &= -y_1 - 2y_1y_2 \\ y_4 &= \dot{y} & y_4' &= -y_2 - y_1^2 + y_2^2 \end{aligned}$$

These four first order equations may now be integrated using the Fortran routine LSODA

Results



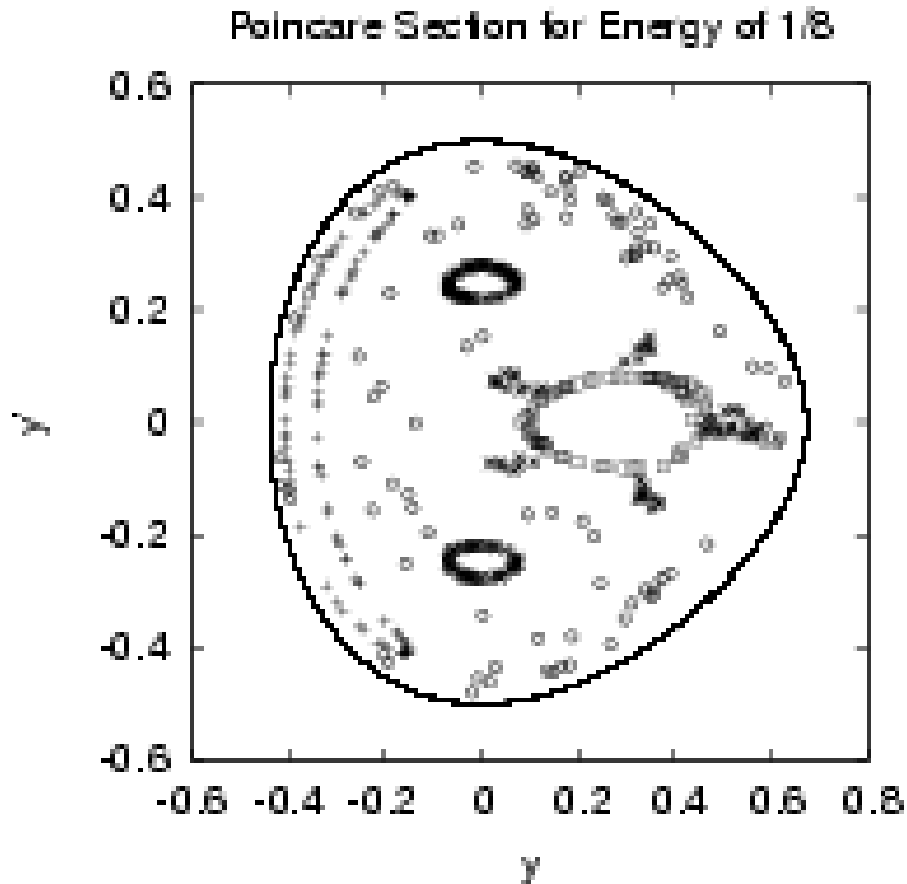
Poincare section for $E=1/12$

Four regions with elliptical orbits

At the center of each region
is an elliptical fixed point

Where the boundary of each
region meets is a hyperbolic
point

If we increase the energy to $E=1/8$ we start to see chaotic behavior

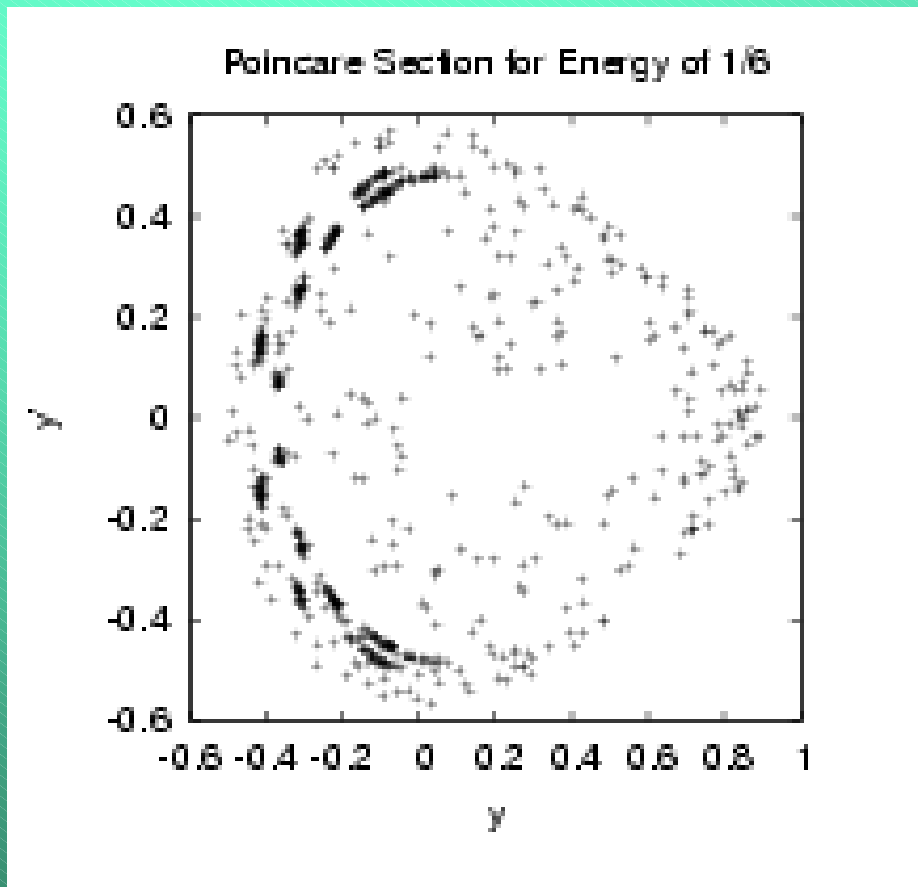


There are still “islands” of non-chaotic behavior with elliptic fixed points at their centers

However, the regions between these islands are now filled with a completely random set of points generated from a single set of initial conditions

These bounded areas of chaos are cross sections of a strange attractor

If the energy is increased to $E=1/6$ the strange attractor pretty much fills up all of the available phase space



This Poincaré map was created with a single set of initial conditions